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**11**study on the method of short-time approximation – Criteria for applicability **Alvin Casandra** <sup>a</sup>, **Suryadi Ismadji** <sup>b</sup>, **Boris A. Noskov** <sup>c</sup>, **Libero Liggieri** <sup>d</sup>, **Shi-Yow Lin**

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**14**article info **Article history: Received 14 November 2014 Received in revised form 25 June 2015 Accepted 6 July 2015 Keywords:**

Surfactant Short-time approximation Diffusivity Dynamic surface tension Ward–Tordai equation Planar surface abstract Despite

**4**its widespread use in the determination of adsorption mechanisms and the estimation of surfac- tant diffusivity,

the short-time approximation method, used for linearly fitting experimental dynamic surface tension data, should be validly applied

**2**only over a very specific range of time intervals or surface pressures.  
Therefore, the definition of general criteria for the applicability of this

method and for error evaluation in diffusivity estimations is fundamental.

**2**In this work, a theoretical numerical simulation of the short-time approximation method was conducted, and general benchmarks for

its accurate utilization were investigated. Specifically, for systems assuming planar gas–liquid surfaces, diffusion-controlled kinetics and a Langmuir adsorption isotherm, simple rules were developed

**2**in terms of limiting surface pressure ( $p_{max}$ ) and dimensionless time ( $t/\tau_{max}$ ) as a function of dimensionless surfactant concentration ( $C_0/a$ ).

For values greater than the limiting (maximal) conditions, the dynamic surface tension curve deviates from the short-time approximation straight line, and thus, the corresponding linear fitting could lead to significant errors in evaluating the diffusivity. The simple criteria proposed in this study thus precisely define the range of applicability for the short-time approximation method. Ó 2015 Elsevier Ltd. All rights reserved. 1.

Introduction Surfactants are essential agents in

**3**sundry practical applications and products, including detergents, inks, adhesives, pesticides, and cosmetics

[1]. In recent years, the global market for surfactants has shown rapid growth, with a 3.8% annual increase: in 2012, the market value was estimated at approximately US\$26.8 billion, and it is expected to reach US\$31.1 billion by 2016 [2]. Because surfactant optimization depends on specific knowledge of their dynamic adsorption behavior [3], it is not surprising that this field of study

**3**has gained tremendous attention over the past two decades.

In 1946, Ward and Tordai introduced a general equation for interpreting the surfactant adsorption kinetics of planar gas–liquid surfaces [4]. In their model, the authors assumed that surfactant diffusion from the bulk to the sublayer is the limiting step compared to surfactant

**10**transfer from the sublayer to the surface. This diffusion-

controlled approach has been observed to be valid for  $\uparrow$

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the majority of small, pure surfactants and surfactant mixture systems [5] such that the equation describing their behavior is considered a fundamental starting point in numerous models used today. However, the Ward–Tordai equation is rather complex, and its application has been hindered for decades by its complicated numerical calculations. Therefore,

**3****in 1959, Defay and Hommelen [6] and Hansen and Wallace [7] introduced a simplified form of the equation, assuming that surfactant backward diffusion**

from the sublayer to the bulk solution could be omitted for initial short time intervals.

**3****In 1979, van den Bogaent and Joos [8] coupled this simplified equation with the** Gibb's adsorption equation and

von Szyszkowski's equation, developing an easy-to-use linear approximation for the Ward–Tordai equation. Due to its simplicity, this short-time linear approximation equation has been widely used to date [9–41]. Specifically, many researchers [10,14–17,19,20,22,23,25,27–32,34,35,37,38,40] have utilized the equation to evaluate surfactant diffusivities or determine the adsorption mechanism of various systems, linearly fitting dynamic surface tension data with the short-time linear approximation equation ( $c$  vs.  $t^{1/2}$ ) for specific ranges of “short” time intervals (generally  $t^{1/2} = 0–5s^{1/2}$ , but up to  $t^{1/2} = 200s^{1/2}$  in [22]). A detailed literature review of these diffusivities, estimated by using the short-time approximation method, is reported in Table 1. Intuitively, because dynamic surface tension data can be accurately linearly fitted only for specific initial time intervals, an appropriate choice of the time range plays a crucial role in the correct evaluation of diffusivity. In fact, it has been reported [22,23,25,28] that quite diverse values of diffusivity have been observed when different time ranges were chosen. Ultimately, thus, certain general criteria for validly applying the short-time approximation method would be essential to guiding researchers in correctly estimating diffusivities. Therefore,

**2****in this work, a theoretical numerical simulation of the short-time approximation method was conducted, and general benchmarks for**

its accurate utilization were investigated. All calculations were carried out for a

16 **diffusion-controlled adsorption** process involving **the mass transport of surfactant molecules from** a uniform **bulk** phase **to**

a freshly created air–water interface. The following conditions were applied: (1) a Langmuir adsorption isotherm, (2) a planar air–water interface, (3) room temperature 25 °C, and (4) a diffusivity of  $D_{set} = 5 \cdot 10^{-6} \text{ cm}^2/\text{s}$  for different dimensionless surfactant concentrations (the

4 **ratio between the bulk concentration and the surfactant activity)**

$C_0/a$ . Additional simulations were conducted for different  $a$  values. 2. Theoretical framework 2.1. Ward–Tordai equation In the case of one-dimensional

7 **diffusion and adsorption onto a planar surface from a bulk phase**

initially containing a uniform distribution of a surface active solute, the diffusion of the surfactant in the bulk continuous phase is described by Fick's law:  $\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial z^2}$ ;  $t > 0$ ;  $z > 0$ ; where  $z$  is the distance from the surface and  $C(z, t)$  is the bulk concentration of surfactant molecules. The boundary and initial conditions for Eq. (1) are as follows:  $C(z=0, t) = C_0$ ;  $\frac{\partial C}{\partial z}(z=0, t) = 0$ ;  $C(z, t=0) = C_0$ ;  $\frac{\partial C}{\partial z}(z, t=0) = 0$ ;  $C(z, t) \rightarrow 0$  as  $z \rightarrow \infty$ . Table 1 A literature review of experimental estimates of the diffusivity ( $D$ ) for different systems, predicted by using the short-time approximation method for different time intervals ( $t$ ). Refs Compound  $D$  ( $10^{-6} \text{ cm}^2/\text{s}$ )  $C_0$  (mol/cm<sup>3</sup>)  $t^{1/2}$  (s<sup>1/2</sup>)  $p$  a (mN/m) Model parameter  $e$   $C_1$

6 **( $10^{-10} \text{ mol}/\text{cm}^2$ ) a ( $10^{-10} \text{ mol}/\text{cm}^3$ )**

13 Glu8-2-Glu8 Glu12-2-Glu12 16 FluoroCarbon4 17 C9-TGE b 19 C12-Maltose ester 20 ANHG550 c 22 Pluronic F68 23 C10E6 25 di-C8 27 C8DMPO d C10DMPO d 28 C10E4 29 b-lactoglobulin b-casein 30 DPPC 31 Hexanol 32 MTAB 33 Triton X-405 34 TDSNa 35 SDS 37 Na-myristate + Na-Decanoate 38, 40 DC10PO DC12AO 1.5 3.5 10 6 0.17 2.3 1 10 8 2 10 8 1.1 2.6 10 8 3.4 0.8 10 6 2.4 0.32 10 6 2.6 1.2 10 6 0.92 0.17 10 11 0.92 0.72 10 11 57 8 10 8 3.0 2 10 7 5.3 4 10 7 20 10 7 4.8 2 10 7 20 10 7 4.9 0.05 10 6 0.6 10 6 4.9 0.1 10 9 5 10 9 4.9 0.1 10 9 5 10 9 12 2 10 9 8 10 10 9 1.5 5 10 6 0.39 10 10 6 13.4 4 10 8 0.7 2.54 10 8 7.63 10 8 2.0 2 10 6 4.7 0.5 10 6 2.8 2 10 6 1.1 2.25 10 7 + 1 10 5 0.029 0.29 10 7 0.011 1.1 10 7 0.4 4.37 10 8 17 10 8 0.1–0.3 6.0 0.1–0.2 0.6 0–14 6.0 0–14 13 0–3 3.5 0–0.3 3.0 0–0.65 2.0 0–0.2 10 0–200 3.0 0–30 4.0 0–0.4 22 0–0.55 7.0 0–0.09 2.0 0–0.05 7.5 0–0.2 1.5 0–0.05 17 0–1.6 2.0 0–0.9 17 0–30 0.5 0–12 17 0–30 1.0 0–30 14 0–19 5.8 0–17 15 0–0.07 9.0 0–0.03 12 0–1 5.9 0–0.21 0.1 0–0.20 1.0 0–0.09 14 0–3 4.0 0–3.1 12 0–1 7.3 0–8 4.4 0–6 11 0–10 7.2 0–8 15 3.99 4.23 4.39 0.31 10.0 63.8 2.46 8.06 3.34 55.8 2.13 4.40 4.84 0.021 2.30 10.8 2.81 1.68 3.20 3830e 3.60 420e 3.18 25.7 1.88 2.02 4.87 0.77 2.79 2.71e 6.86 52100 2.80 3600e 0.92 0.22 9.24 4030e 3.84 16000e 3.08 38800e 3.70 330e 4.20 280e a Surface pressure,  $p = c_0 c$ . b TGE = tryptophan glycerol ether surfactant. c ANHG550 = heterogemini surfactant of PEG with MW = 550. d DMPO = dimethyl phosphine oxides. e Parameters ( $C_1$  and  $a$ ; maximum surface concentration and surfactant activity) obtained from previous studies. where  $C_0$  is the initial bulk

concentration,  $C$  is the surface concentration of surfactant and zero  $C$  indicates an initially clean air–water surface. By using the Laplace transform, the adsorption of surfactant molecules

12 **as a function of time can be** formulated **in terms of**

an unknown sublayer concentration  $C_S(t) = C(z = 0, t)$ :  $C_0 \frac{1}{4} 2C_0 \frac{Dt}{1+2D} \frac{1}{p} \frac{1}{p} \frac{1}{2} Z \frac{p}{\text{fit}} C_S \delta t \frac{s}{p} d$   
 $\frac{1}{\text{ffiffiffi}} \delta 6 \frac{1}{p} 0$  where  $C_S$

5 **is the** sublayer **concentration and D is the surfactant** diffu- sivity. **This relation is**  
**the**

well-known Ward–Tordai equation, which is applicable for an initially clean planar gas–liquid surface and a diffusion-controlled surfactant mass transport process. This sublayer concentration can be determined numerically [42] by a modification of the technique used by Miller and Kretzschmar [43]. Specifically, the first and second terms on the right-hand side of Eq. (6), defined in this work as TermF and TermB, are linked to surfactant forward

10 **diffusion from the bulk phase to the sublayer and** surfactant backward  
diffusion **from the** sub-layer **to the**

bulk solu- tion, respectively. 2.2. Dynamic surface tension To obtain the theoretical dynamic surface tension curve, the Ward–Tordai equation (6) has to be combined with two additional relations: the

13 **adsorption isotherm and the equation of state. The** adsorption isotherm  
relates **the surface**

concentration and subsurface concentration of surfactant molecules at constant tem- perature. When

15 **it is assumed that the mass transport** process **is** diffusion **-controlled** and **the**

adsorbed surfactants on the surface do not interact with each other, the Langmuir isotherm can describe the

13 **relationship between the surface** concentration **and** the **bulk surfactant**  
**concentration:**  $C_1 \frac{1}{4} \times \frac{1}{4} C$

$C \frac{1}{p} C \delta 7 \frac{1}{p}$  where  $C_1$  and  $a$  are the model parameters of the Langmuir iso- therm, and  $x$  is the dimensionless surface concentration. Parameter  $C_1$



approximation technique for identifying adsorption mechanisms and estimating diffusivity, a theoretical numerical simulation was performed in this study. All calculations were carried out for an established diffusivity of  $D_{set} = 5 \times 10^{-6} \text{ cm}^2/\text{s}$  and for different dimensionless bulk concentrations  $C_0/a$ . Initially, by using the Ward–Tordai equation and the Langmuir adsorption isotherm, with parameters  $C_1 = 5$

**$6 \times 10^{10} \text{ mol}/\text{cm}^2$  and  $a = 1 \times 10^9 \text{ mol}/\text{cm}^3$**

as delineated in Section 2, the exact relaxation curves of the dynamic surface tension,  $c(t)$ , and the surface concentration,  $C$ , were generated for five different values of  $C_0/a$ , as shown in Fig. 1a. Plotting the surface tension (or surface pressure) as a function of  $t^{1/2}$ , initially linear trends (as predicted by Eq. (12)) were clearly identified for specific short time intervals (Fig. 1b). A decrease in the time intervals (denoted  $t_{max}$  in this study) is noted, in which the theoretical curves show linear behavior with increasing bulk concentration for  $a = 1 \times 10^9 \text{ mol}/\text{cm}^3$ . The data in Fig. 1b indicate that  $t_{1/2}^{max}$  decreases from  $80 \text{ s}^{1/2}$  to  $32 \text{ s}^{1/2}$  as  $C_0/a$  increases from 0.4 to 8. In other words,  $t_{max}$  decreases from 6400 s to 1024 s as  $C_0/a$  increases from 0.4 to 8. A surface pressure (denoted  $p_{max}$ ) corresponding to  $t_{max}$  is therefore defined, in which the theoretical curves show a linear  $c$  vs.  $t^{1/2}$  behavior. Fig. 2 shows the change in  $t_{max}$

**5 as a function of surfactant concentration for surfactants with different**

surfactant activities,  $a = 1, 5, \text{ and } 20$  ( $10^9 \text{ mol}/\text{cm}^3$ ). For surfactants with different activities,  $t_{max}$  always decreases with bulk concentration. Moreover, a dramatic decrease occurs at a  $C_0/a$  value of approximately 8. Fig. 3 details the deviation in  $c(t)$  between the tension data ( $c_{data}$ ) predicted by the Ward–Tordai equation and those ( $c_{approx}$ ) determined by the short-time approximation (Eq. (12)) at  $t < t_{max}$ ;  $D_c(t) = c_{data} - c_{approx}$ . The data in Fig. 3 indicates for the case of  $C_0/a = 8$ , a positive deviation occurs at the beginning, followed by a negative  $D_c(t)$ ; then, another positive value of  $D_c(t)$  occurs. The maximum deviation in  $c$  in this case is  $0.7 \text{ mN}/\text{m}$ , which is roughly the measurement uncertainty in our experiments. Fig. 4 summarizes the relaxation of  $D_c(t)$  for different surfactant concentrations,  $C_0/a = 0.4\text{--}30$ . At intermediate surfactant concentrations (2.2–8), the deviation in  $D_c(t)$  is similar to the behavior shown in Fig. 3. When the concentration is low ( $C_0/a = 0.4\text{--}2$ ),  $D_c(t)$  shows a small negative value, followed by a small positive one. In a b Fig. 1. Relaxation profiles generated using the Ward–Tordai equation and the Langmuir isotherm ( $C_1 = 5$

**$6 \times 10^{10} \text{ mol}/\text{cm}^2, a = 1 \times 10^9 \text{ mol}/\text{cm}^3,$**

and  $D_{set} = 5 \times 10^{-6} \text{ cm}^2/\text{s}$ ). Profiles of (a)  $c(t)$ ,  $C(t)$  and (b)  $p(t^{1/2})$  for five different values of  $C_0/a$  ( $=0.4, 1, 2, 4, 8$ ). 104 A 103 t max (s) 102 B 10 C 1 10-1 0.1 1  $C_0/a$  10 Fig. 2. Relaxation of  $t_{max}$

**2 as a function of surfactant concentration  $C_0/a$**

and surfactant activity:  $a = 1$  (A), 5 (B), and 20 (C) ( $10^9 \text{ mol}/\text{cm}^3$ ). 70  $\gamma$  (mN/m) 60 50 0.8 0.4 0  $\Delta\gamma$  (mN/m)  $t^{1/2}_{max} = 32 -0.4 C_0/a = 8 -0.8 0 10 20 30 40 t^{1/2}$  (s $^{1/2}$ ) Fig. 3. Deviation in surface tension  $D_c(t)$  between the  $c(t)$  data (solid curve,  $c_{data}$ ) and the linear prediction (dashed line,  $c_{approx}$ ) of the short-time

approximation for  $C_0/a = 8$ ;  $D_c(t) = c_{data} - c_{approx}$ . 0.8  $C_0/a=8$  6 0.4 4 3.5 2.5 3  $\Delta\gamma$  (mN/m) 2 0.4 0 1 4 0.1 6  $C_0/a=3.5$  3 2.2 2.2 1 -0.4 8  $\Delta\gamma$  2.5 0 2 2.5 2 3 3.5 0.4 30 12 10 -0.8 0 10 20  $t^{1/2}$  30 40 50 0 20 40  $t^{1/2}$  (s $^{1/2}$ ) 60 80 Fig. 4. Relaxation of the tension deviation  $D_c(t)$  at different surfactant concentrations  $C_0/a$ . contrast, at high concentrations ( $C_0/a > 10$ ), a significantly positive value of  $D_c(t)$  occurs first, followed by a negative one. In applying the short-time approximation technique, experimental dynamic surface tension data are linearly fitted (lines 1 in Fig. 5) to estimate the surfactant diffusivity when an adsorption process is diffusion-controlled. Fig. 5 illustrates the fitting and the evaluation of the diffusivity  $D$  using the short-time approximation technique for the early stages of the dynamic  $c(t)$  data derived from the Ward–Tordai equation. For surfactant concentrations ( $C_0/a$ ) lower than 2 this linear fitting calculation leads to an underestimation of the diffusivity (Fig. 5a), whereas for  $C_0/a$  values greater than 2,  $D$  is overestimated (Fig. 5c). In both cases, the deviation from the established value  $D_{set}$  increases with the time range considered such that, intuitively, shorter time intervals lead to more precise estimates of the diffusivity (lines 2 in Figs. 5a and c). Interestingly, as previously noted, for  $C_0/a = 2$ , the deviations in the estimated  $D$  (shown as the line 1 in Fig. 5b) from  $D_{set}$  (shown as the line “2” in Fig. 5b) are negligible for all time ranges, where the  $c(t) t^{1/2}$  curve follows a linear trend. This deviation in the estimated value of  $D$  (underestimation at low  $C_0/a$ , nearly zero deviation at  $C_0/a = 2$ , and overestimation at large  $C_0/a$ ) is derived from the deviation in the tension  $D_c(t)$  shown in Fig. 4. At low concentration,  $D_c(t) < 0$  first, then  $D_c(t) > 0$  over time (Figs. 4 and 5a). The positive  $D_c(t)$  deviation causes a smaller  $D$  value to be estimated by linear fitting (Eq. (12)); therefore,  $D$  is underestimated. At high concentration,  $D_c(t) > 0$  first, then  $D_c(t) < 0$  over time (Figs. 4 and 5c). The negative  $D_c(t)$  deviation causes a larger  $D$  value to be estimated by linear fitting; therefore a  $D$  is overestimated. At  $C_0/a = 2$ ,  $D_c(t)$  is small; therefore, the effects induced by positive and negative values of  $D_c(t)$  are balanced, resulting in a small deviation in  $D$ . The data in Figs. 5a and c lead to another conclusion: a larger deviation in  $D$  is observed when a larger range of  $c(t)$  data is used for the short-time approximation. Fig. 5c shows that  $D = 7.7 \cdot 10^{-6}$  cm $^2$ /s was obtained when the  $c(t)$  data at  $t^{1/2} < 20$  s $^{1/2}$  were used. When the  $c(t)$  data at  $t^{1/2} < 32$  s $^{1/2}$  were used, a larger deviation in  $D$  ( $8.6 \cdot 10^{-6}$  cm $^2$ /s) was observed. The effect of surfactant concentration ( $C_0/a$ ) on the deviation in  $D$  ( $D_{err} = 100(D_{approx} - D_{set})/D_{set}$ ) is shown in Fig. 6. A nearly linear dependence was observed when  $D_{err}$  was plotted

17as a function of surface pressure  $p$ . The value of

$p$  indicates the range of  $c(t)$  data used for linear fitting by the short-time approximation technique. This linear dependence confirms the conclusion discussed above: a larger deviation in  $D$  is observed when a larger range of  $c(t)$  data is used. Fig. 6a also indicates that at high concentrations, a larger range of  $c(t)$  data can be used for evaluating  $D$  in applying the short-time approximation technique. The last data point (solid circle) indicates the maximum range of  $c(t)$  data over which the tension deviation [ $D_c(t) = c_{data} - c_{approx}$ ] is still allowable (less than the measurement uncertainty; 0.1 mN/m at low surfactant  $72 a \gamma$  (mN/m)  $D = 4.07$  (10 $^{-6}$  cm $^2$ /s)  $71 2 3.38 3 2 1 70 1 0 20 40 60 t^{1/2}$  (s $^{1/2}$ )  $100 72 b 70 \gamma$  (mN/m)  $68 D = 4.94$  (10 $^{-6}$  cm $^2$ /s)  $66 64 1 2 1 0 10 20 30 40 t^{1/2}$  (s $^{1/2}$ )  $60 70 c D = 7.68$  (10 $^{-6}$  cm $^2$ /s)  $\gamma$  (mN/m)  $60 8.63 2 2 3 50 1 1 0 10 20 t^{1/2}$  (s $^{1/2}$ )  $30 40 50$  Fig. 5. The linear behavior,  $c(t) \cdot t^{1/2}$ , determined by the short-time approximation at different surfactant concentrations for the data shown in Fig. 1 (line 1):  $C_0/a = 0.4$  (a), 2.0 (b), and 8.0 (c). Lines 1 and 2 shows the effect of the time interval used on estimating the diffusivity (lines 2 use a short interval than lines 1). Line 3 shows  $c(t)$  derived from the short-time approximation with  $D_{set} = 5 \cdot 10^{-6}$  cm $^2$ /s.  $80 a C_0/a = 8$   $21.0 40 D_{err}$  (%)  $4 12.6 0 2 1 0\%$   $0.4 6.22 3.44 -40 1.66 0 6 12 18 \pi$  (mN/m)  $80 b C_0/a = 8$   $1.66 40 D_{err}$  (%)  $0.29 -40 0.4 0.25 0 2 0\%$   $1 0.45 4 1.01 0 0.5 t^* 1 1.5$  Fig. 6. Nondimensional diffusivity error ( $D_{err}$ ) as a function of (a) the surface pressure ( $p$ ) and (b) dimensionless time ( $t^*$ ) for  $C_0/a = 0.4, 1, 2, 4, \text{ and } 8$ .

concentrations and 0.7 mN/m at high concentrations in this study). The relationship between Derr and dimensionless time,  $\tau = t/[C2eq/(C20D)]$ , at different surfactant concentrations is shown in Fig. 6b, in which Derr levels off as  $\tau$  increases at high concentrations. To determine these general benchmarks, the surfactant diffusivity, estimated by the short-time approximation technique for different surfactant concentrations, was initially investigated as a function of the surface pressure,  $p = c_0 - c$ . Fig. 7a shows the limiting conditions ( $p_{max}$ ) for which the deviation between  $c_{approx}$  (from the short time approximation straight line) and  $c_{data}(t)$  (from the Ward–Tordai equation) is allowable ( $D_c(t) < 0.1$  and 0.7 mN/m at low and high concentrations). Fig. 7a shows the region (below the curve of  $p_{max}$  vs.  $C_0/a$ ) where the short-time approximation is applicable (i.e., with a reasonable  $D_c(t)$ ). The error in estimating the surfactant diffusivity may range from 30% to 70% (Fig. 6), depending on the surfactant concentration  $C_0/a$  and the region of  $c_{data}(t)$  used (indicating by  $p$  or  $\tau$ ). In order for the readers to apply the data (the region where the short-time approximation is applicable) easily, Fig. 7a are also shows this dependence in dimensionless parameters  $p^*_{max} (=p_{max} c_0 / c_{eq})$  and  $\tau^*_{max} (=t_{max}/[C2eq/(C20D)])$

### 2as a function of surfactant concentration $C_0/a$ .

Dimensionless time  $\tau^*_{max}$  indicates the time  $\tau$  where  $p$  reaches  $p^*_{max}$ . Interestingly, as illustrated in Fig. 7b, a linear dependence ( $p_{max}$  vs.  $C_0/a$ ) was observed at low surfactant concentrations ( $C_0/a \leq 6$ ). The existence of a linear dependence is probably because the surfactant concentration is low. It is noted that this linear dependence in term of dimensionless parameter ( $p^*_{max}$  vs.  $C_0/a$  and  $\tau^*_{max}$  vs.  $C_0/a$ ) works only at  $0.4 \leq C_0/a \leq 6$ . At an extremely low concentration ( $C_0/a \leq 0.4$ ), both  $p^*_{max}$  and  $\tau^*_{max}$  decrease as concentration increases. a 20  $\pi^*_{max}$   $\pi_{max}$  (mN/m) 10  $\pi_{max}$  0  $t^*_{max}$  0.01 0.1  $C_0/a$  1 10  $\pi_{max}$  (mN/m) 10 20 0.8 0.6  $\pi^*_{max}$  0.4 0  $C_0/a$  4 20  $\pi_{max}$  10 8 1.8 1  $t^*_{max}$  0.2 b 0 0 2  $C_0/a$  4 0 0 4  $C_0/a$  8 6 2.6 1.8  $t^*_{max}$  1 0.2 1.8  $t^*_{max}$  1 0.2 0.8 0.4  $\pi^*_{max}$  0 0.8  $\pi^*_{max}$  0.6 0.4 Fig. 7. (a) Conditions (maximal surface pressure  $p_{max}$ , dimensionless time  $\tau^*_{max}$ , and dimensionless surface pressure  $p^*_{max}$ ) under which the short-time linear approximation is applicable; (b) conditions for surfactant concentration  $C_0/a \leq 6$ . A simple linear dependence can be used to determine the limiting working conditions of the short-time approximation:  $p^*_{max} = 1/4 \cdot 2.80 \delta C_0 = a \cdot p$   $p = 0.72 \cdot 1/2 \cdot 1/4 \cdot \delta \text{mN} = m \cdot \delta 13 \cdot p$   $t^*_{max} = 1/4 \cdot 0.21 \delta C_0 = a \cdot p$   $p = 0.12 \cdot 1/2 \cdot 1/4 \cdot \delta 14 \cdot p$   $p^*_{max} = 1/4 \cdot 0.066 \delta C_0 = a \cdot p$   $p = 0.33 \cdot 1/2 \cdot 1/4 \cdot \delta 15 \cdot p$  For surface pressures greater than this maximal surface pressure,  $p > p_{max}$ , the dynamic surface tension curve deviates significantly from the linear region of the short-time approximation, and thus, the linear fitting of the tension could lead to significant errors in the calculated diffusivity. Therefore, Eq. (13) defines the range

### 16of applicability of the short-time

linear approximation technique

### 12in terms of the surface pressure as a function of surfactant concentration

$C_0/a$ . A further analysis was conducted

18to evaluate the possible effect of the surfactant activity on the

nondimensional error in the diffusivity evaluated from the short-time approximation technique. As illustrated in Fig. 8, the Derr-c curves for different a values ( $a = 1 \cdot 10^{-9}$ ,  $5 \cdot 10^{-9}$  and  $20 \cdot 10^{-9}$  mol/cm<sup>3</sup>), but the same surfactant concentration  $C_0/a$ , present a distinct trend. It can therefore be concluded that the surfactant activity plays only a minor role in obtaining a valid diffusivity estimate. 4. Discussion Based on the definition of the Ward-Tordai equation in Eq. (6) and the assumption that Eq. (9) is valid when surfactant backward diffusion can be assumed to be negligible, it is trivial to show how the deviation of c from the linear fitting is related to the actual 0 0% Derr (%) -10 a  $C_0/a=0.4$  -20 -30 72 71  $\gamma$  70 Derr (%) 0 -10 b  $C_0/a=2$  0% 10 40 Derr (%) 20 c  $C_0/a=6$  60 7 2 70  $\gamma$  68 66 0 0% 7 2 66  $\gamma$  60 54 Fig. 8. A nearly unique nondimensional error in diffusivity Derr evaluated from the short-time approximation technique as function of surface tension [c (t) range used for the linear short-time approximation fitting] for different surfactant activities  $a = 1$  (s), 5 (r), and 20 ( ) ( $10^{-9}$  mol/cm<sup>3</sup>) at surfactant concentrations of  $C_0/a = 0.4$  (a), 2 (b), and 6 (c). balance between the forward and backward surfactant mass transport processes. To enhance understanding of this crucial aspect, both TermF ( $=2C_0(Dt/p)^{1/2}$ , Fig. 9a, solid curves) and TermB ( $=2(D/p)^{1/2} R_{ptffi} CS\dot{t} s\beta dpffisffiffi$ ; Fig. 9a, dashed curves) of the 0 Ward-Tordai equation were calculated, using the Langmuir isotherm ( $C_1 = 5$

610 10 mol/cm<sup>2</sup>,  $a = 1 \cdot 10^{-9}$  mol/cm<sup>3</sup>,

and  $D_{set} = 5 \cdot 10^{-6}$  cm<sup>2</sup>/s), for three different surfactant concentrations. As clearly shown in Fig. 9a, after a rather short initial time interval, the backward diffusivity can no longer be omitted. Indeed, it is worth noting that even for time ranges close to  $t = 0$ , TermB is not null and a small degree of backward diffusion occurs. The relative importance of the backward diffusivity compared to the forward diffusivity is also clearly illustrated in Fig. 9b, in which the TermB/TermF ratio is analyzed: for time intervals  $t > 100$ s the diffusion of the surfactant from the sublayer to the bulk phase must be considered, and the use of the short-time linear approximation technique should not be considered to be completely accurate. Finally, we compared all of the diffusivity data, estimated using the short-time linear approximation technique and reported in the literature, with the limiting criteria identified in this study. Fig. 10 shows the surface pressure p as function of surfactant concentration  $C_0/a$  for the data presented in Table 1. The maximum error limiting criteria (solid curve) and the criteria indicating a 30% error in diffusivity (dashed curve) were analyzed. It is noteworthy that the majority of the reported diffusivity data appear to be included Termi ( $10^{-10}$  mol/cm<sup>2</sup>) 8 6 4 0 Termi ( $10^{-10}$ ) 6 0 0 2 a  $C_0/a = 8$  F 8 B 3 0.4  $t^*$  2 F 2 B 6  $C_0/a = 8$  F 8 B 2 F 2B 0.4 10 12 1 10 102 t (s) 103 0.6  $C_0/a=0.4$   $C_0/a = 8$  2 TermB / TermF 0.6 TermB / TermF 8 0.3 0.4 0 0 3  $t^*$  6 2 0.2 b 0.4 0 1 10 102 t (s) 103 Fig. 9. (a) Forward diffusion (solid curves) and backward diffusion (dashed curves) determined from the Ward-Tordai equation and (b) the ratio of backward diffusion (TermB) to forward diffusion (TermF)

18as a function of time (main figure) and dimensionless time (inset);  $t/$

$t/ = 25,500$ , 5560, and 617 for  $C_0/a = 0.4$ , 2.0 and 8.0, respectively. in our error limiting criteria, with the following error range for the estimated diffusivity coefficient:  $30\% < Derr < 30\%$ . Regarding the various data reported in the literature that were observed to fall above our maximum limiting criteria, special attention is

required. Specifically, we do not claim that the surfactant diffusivities estimated by the authors who furnished these data are actually affected by errors to a great extent, but the fact that these points are far above our limit certainly suggests that the fundamental assumptions that led to our criteria (planar gas–liquid surfaces, diffusion-controlled mass transport process and Langmuir adsorption isotherm) may not be applied for these specific systems. In this work, a fixed diffusivity ( $D_{set} = 5 \cdot 10^{-6} \text{ cm}^2/\text{s}$ ) was set for all the theoretical simulation. This diffusivity was chosen because the most of small surfactants have a diffusivity closed to this value. It is noted that the choice of the  $D_{set}$  won't make any change on the conclusion and phenomenon presented in this manuscript because the data was presented in the variation percentage (Derr), as shown in Figs. 6 and 8. This study considered only a planar air–water interface for surfactants following the Langmuir adsorption isotherm (i.e., no intermolecular

**15 interaction between the adsorbed surfactant molecules** at surface). However, **the**

evaluation error using the short time approximation method depends not only on surfactant concentration ( $C_0/a$ ), but also on the spherical curvature of air–water  $20 \cdot 15 \cdot \pi$  (mN/m)  $10 \cdot 5 \cdot 0 \cdot 73\% \cdot 23 \cdot 51\% \cdot 27 \cdot 28 \cdot 29 \cdot 38 \cdot 30 \cdot 35 \cdot 31 \cdot 1368 \cdot 264 \cdot 7768\% \cdot 80\% \cdot 85\% \cdot 35 \cdot 27 \cdot 31 \cdot 88\% \cdot 29 \cdot 35 \cdot 5\% \cdot 27 \cdot 38 \cdot 37 \cdot 30 \cdot 28 \cdot 32 \cdot 16 \cdot 38 \cdot -1\% \cdot -44\% \cdot -32\% \cdot 22 \cdot 217' \cdot 35 \cdot 22 \cdot 30\% \cdot 27 \cdot 30 \cdot 30 \cdot 28 \cdot 19 \cdot 27 \cdot 27 \cdot 1278 \cdot 1177 \cdot 15 \cdot 19 \cdot 27 \cdot 29 \cdot 19 \cdot 29 \cdot 0.1 \cdot 1 \cdot C_0/a \cdot 10 \cdot 100$  Fig. 10. Comparison between literature data (in Table 1) and limiting surface pressure  $p$  criteria proposed in this work as function of  $C_0/a$ . The criteria are evaluated for maximum limiting criteria (solid curve) and for  $30\% < Derr < 30\%$  (dashed curve). Literature data shown in circles were obtained by pendant drop or maximum bubble pressure measurement methods, data shown in squares were obtained by the Wilhelmy plate method, and data shown in diamonds were obtained by the short-time linear approximation technique but without passing through the initial point of  $c_0 = 72 \text{ mN/m}$  at  $t = 0$ . The numbers shown near (or inside) the symbols indicate the reference articles. The numbers shown near the line indicate the theoretical percentage of diffusivity error. interface and the intermolecular interaction force. The applicability criteria for these more complicated systems have been studied and in the course of the article writing in our laboratory for the time being. 5. Conclusions The short-time approximation technique is widely used for linearly fitting experimental dynamic surface tension data to estimate the diffusivity and to determine the adsorption mechanism of a given process. However, because it is assumed that at the beginning of the adsorption process, surfactant backward diffusion from the sublayer to the bulk phase is negligible and  $\Gamma$  is constant, the applicability of this technique is limited to specific time intervals. Therefore,

**2 in this work, a theoretical numerical simulation of the short-time approximation technique was conducted,** and general criteria for

its accurate utilization were analyzed. Specifically, for systems assuming planar gas–liquid surfaces, a diffusion-controlled mass transport process and a Langmuir adsorption isotherm, some simple rules were established in terms of maximal

**4 surface pressure  $p_{max}$  and maximal dimensionless time  $t_{m/ax}$  as a function of the dimensionless surfactant concentration  $C_0/a$ .**

For  $p > p_{\max}$  or  $t > t_{\max}$ , the dynamic surface tension curve deviates from the linear region of the short-time approximation, and thus, the linear fitting of the surface tension could lead to significant errors in diffusivity estimates. Under these conditions, thus, the actual applicability of the short-time linear approximation method must be critically evaluated. Conflict of interest None declared. Acknowledgments The authors would like to express great appreciation to Taiwan NSC and the Italian National Research Council for the support within the project "Mixed Particle-Surfactant Layers at Liquid Interfaces" (NSC/CNR Bilateral Agreement 2014-2015 and NSC-100-2221-E-011-101). References [1] B.E. Chistyakov, Theory and practical application aspects of surfactants, *Surfactants Chem. Interfacial Prop. Appl.* (2001) 511–618. [2] Acmite Market Intelligence, Market Report: Global Surfactant Market, 2013. [3] S.S. Dukhin, R. Miller, G. Kretzschmar, On the theory of adsorption kinetics of ionic surfactants at fluid interfaces, *Colloid Polym. Sci.* 269 (1991) 923–928. [4] A.F.H. Ward, L. Tordai, Time-dependence of boundary tensions of solutions I. the role of diffusion in time-effects, *J. Chem. Phys.* 14 (1946) 453–461. [5] R. Miller, A.V. Makievski, V.B. Fainerman, Dynamics of adsorption from solutions, *Surfactants Chem. Interfacial Prop. Appl.* (2001) 287–399. [6] R. Defay, J.R. Hommelen III, The importance of diffusion in the adsorption process of some alcohols and acids in dilute aqueous solutions, *J. Colloid Sci.* 14 (1959) 411–418. [7] R.S. Hansen, T.C. Wallace, The kinetics of adsorption of organic acids at the water-air interface, *J. Phys. Chem.* 63 (1959) 1085–1092. [8] R. van den Bogaert, P. Joos, Dynamic surface tensions of sodium myristate solutions, *J. Phys. Chem.* 83 (1979) 2244–2248. [9] C.D. Ampatzidis, E.-M.A. Varka, T.D. Karapantsios, Dynamic surface properties of eco-friendly phenylalanine glycerol ether surfactants at the W/A interface, *Colloids Surf. A Physicochem. Eng. Asp.* 441 (2014) 872–879. [10] M. Eftekhardakhah, P. Reynders, G. Øye, Dynamic adsorption of water soluble crude oil components at air bubbles, *Chem. Eng. Sci.* 101 (2013) 359–365. [11] S. Ferdous, M.A. Ioannidis, D. Henneke, Adsorption kinetics of alkanethiol- capped gold nanoparticles at the hexane–water interface, *J. Nanoparticle Res.* 13 (2011) 6579–6589. [12] T.D. Gurkov, Adsorption kinetics under the influence of barriers at the subsurface layer, *Colloid Polym. Sci.* 289 (2011) 1905–1915. [13] K. Sakai, S. Umezawa, M. Tamura, Y. Takamatsu, K. Tsuchiya, K. Torigoe, et al., Adsorption and micellization behavior of novel gluconamide-type gemini surfactants, *J. Colloid Interface Sci.* 318 (2008) 440–448. [14] A.M. Díez-Pascual, A. Compostizo, A. Crespo-Colín, R.G. Rubio, R. Miller, Adsorption of water-soluble polymers with surfactant character. Adsorption kinetics and equilibrium properties, *J. Colloid Interface Sci.* 307 (2007) 398–404. [15] T. Yoshimura, A. Sakato, K. Tsuchiya, T. Ohkubo, H. Sakai, M. Abe, Adsorption and aggregation properties of amino acid-based N-alkyl cysteine monomeric and N, N'-dialkyl cysteine gemini surfactants, *J. Colloid Interface Sci.* 308 (2007) 466–473. [16] Y.A. Gao, W.G. Hou, Z.N. Wang, G.Z. Li, B.X. Han, G.Y. Zhang, Dynamic surface tensions of fluorosurfactant solutions, *Chinese J. Chem.* 23 (2005) 362–366. [17] C. Delgado, M.D. Merchán, M.M. Velázquez, S. Pegiadou, L. Pérez, M. Rosa Infante, The adsorption kinetics of 1-N-I-tryptophan-glycerol-ether surfactants at the air-liquid interface. effect of surfactant concentration and alkyl chain length, *Colloids Surfaces A Physicochem. Eng. Asp.* 233 (2004) 137–144. [18] J.M. Rodríguez Patino, C. Carrera Sanchez, S.E. Molina Ortiz, R. Rodríguez Nino, C. Anon, Adsorption of soy globulin films at the air–water interface, *Ind. Eng. Chem. Res.* 43 (2004) 1681–1689. [19] U.R.M. Kjellin, J. Reimer, P. Hansson, An investigation of dynamic surface tension, critical micelle concentration, and aggregation number of three nonionic surfactants using NMR, time-resolved fluorescence quenching, and maximum bubble pressure tensiometry, *J. Colloid Interface Sci.* 262 (2003) 506–515. [20] E. Alami, K. Holmberg, J. Eastoe, Adsorption properties of novel gemini surfactants with nonidentical head groups, *J. Colloid Interface Sci.* 247 (2002) 447–455. [21] R. Rodríguez Nino, J.M. Rodríguez Patino, Effect of the aqueous phase composition on the adsorption of bovine serum albumin to the air–water interface, *Ind. Eng. Chem. Res.* 41 (2002) 1489–1495. [22] M.G. Muñoz, F. Monroy, F. Ortega, R.G. Rubio, D. Langevin, Monolayers of symmetric triblock copolymers at the air–water interface. 2. adsorption kinetics, *Langmuir* 16 (2000) 1094–1101. [23] B.V. Zhmud, F. Tiberg, J. Kizling, Dynamic surface tension in concentrated solutions

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